

Photonic molecules modes in coupled spherical microcavities with CdTe nanocrystals

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ABSTRACT

We present a detailed study of the localized coupled-cavity modes in a photonic molecule formed from two dielectric spherical microcavities with CdTe nanocrystals. A layer-by-layer deposition technique provides controllable coating of the microspheres with a shell of close-packed nanocrystals of approximately 4 nm in diameter. The observed spectral structures originate from the coupling of the electronic transitions in nanocrystals and the photon states of interacting microspheres and, in analogy to the formation of molecular electronic orbits, can be assigned to bonding and antibonding photon localized states.

Keywords: coupled microcavities, semiconductor nanocrystals, whispering gallery modes, mode degeneracy, optical delay

1. INTRODUCTION

There is currently great interest in design and fabrication of optical structures for manipulation of light on micrometer-length scale. Much of this attention has been focused on photonic band-gap structures or two-dimensional photonic wires due to the possibility of optical interconnection¹. Closely related to these developments is the intensive work on the modification of photon modes in spherical microcavities, which has been of great interest both for studies of fundamental optical properties and for the potential applications².

Extending the ideas of the linear combination of atomic orbitals method to the classical wave case, it was recently suggested that Mie resonances of single spherical microcavity play the same role as the atomic orbitals in the electronic case and the spatial distributions of whispering gallery modes (WGM) can be described by analogy with the orbitals in a hydrogen atom³. In the absence of gain, the WGM resonances can be characterized by a mode number n (angular quantum number), mode order l (radial quantum number), and azimuthal mode number m (azimuthal quantum number). The value of n is proportional to the circumference divided by the wavelength of the light propagating within the microsphere, the mode order l indicates the number of maxima in the radial distribution of the internal electric field, and the azimuthal mode number m gives the orientation of the WGMs orbital plane. WGMs in a single sphere are degenerate with respect to m because of the spherical symmetry.

The electro-magnetic fields confined in microsphere are given by⁴

$$\mathbf{E}_{TE} = j_n(nkr)\mathbf{X}_{nm}(\theta, \phi), \quad (1)$$

for modes having no radial components of the electric field (transverse electric or TE modes) and

$$\mathbf{E}_{TM} = \nabla \times j_n(nkr)\mathbf{X}_{nm}(\theta, \phi), \quad (2)$$

for the transverse magnetic (TM) modes (no radial component of the magnetic field).

Here $X_{nm}(\theta, \phi)$ are the vector spherical harmonics given in spherical polar, $j_n(n, kr)$ represents the spherical Bessel function, where $k = 2\pi/\lambda$ is the wavenumber, n_r is the refractive index of the sphere and λ is the free space wavelength.

On the other hand, the wave functions Ψ_{slm} for the electron confined in the hydrogen atom are given by⁵

$$\Psi_{slm} = R_{sl}(r)Y_{lm}(\theta, \phi), \quad (3)$$

where $R_{sl}(r)$ is known as the Laguerre polynomials. The eigenfunctions (1) and (3) are very similar and their spatial distributions are characterized by the three integer s , l and m (for hydrogen atom) and n , l and m (for microsphere), which correspond to total angular, radial and the azimuthal quantum numbers, respectively. Based on all above similarities this approach has enabled small dielectric spheres to be considered as “photonic atoms”^{3,6}. However, unlike energy states of electron in the atom, photonic states in spherical microcavities are not localized due to finite storage time (τ) of photons in the resonant mode. This “photon lifetime” is controlled by quality factor of WGM Q and therefore can be limited by diffractive losses, absorption, gain, shape deformation or refractive index inhomogeneity⁴. As a result, the resonant internal field of a spherical cavity is not completely confined to the interior of the microparticle. Depending on the size of the microsphere, the evanescent field can extend into the surroundings up to a couple of micrometers. It was recently recognized that the partial delocalization of Mie resonance states is of great importance because it suggests a possibility for coherent coupling between WGMs of two adjacent spherical particles with closely matched sizes. Such a system of coherently coupled photonic atoms may be called a “photonic molecule” (PM)⁷ and can be employed in order to manipulate photons in the micrometer length scale. In analogy to the formation of molecular electronic orbits, the tight binding approximation provides two combinations for the electromagnetic field in a system of interacting microspheres: bonding (BN) and antibonding (ABN) states⁷⁻¹⁰. Experimentally, the coupling of the photon modes of individual microspheres in the PM can cause a narrow resonance of a photonic atom to split into two modes of lower Q -factor⁸. This phenomenon has been demonstrated in a system of two square, photonic dots coupled by a narrow channel⁷ in a dye-stained bisphere system^{9,11,12}, in photonic dots doped by CdSe nanocrystals (NCs)¹³ and in chains of polymer-blend microparticles¹⁴. However recent theoretical considerations¹⁰ and experimental studies^{15,16} reveal complex internal distribution of density of photonic states of PM originating from lifting of degeneracy of PM modes with respect to the azimuthal index.

2. EXPERIMENT

In this paper we present a detailed study of the localized coupled-cavity modes in a PM formed from two dielectric spherical microcavities with CdTe NCs. Aqueous dispersions of melamine-formaldehyde (MF) microspheres, of $3 \pm 0.05 \mu\text{m}$ in diameter, were combined with luminescent CdTe NCs using a layer-by-layer deposition technique¹⁷. The colloidal solution of CdTe NCs, with a PL maximum at 563 nm (2.9 nm radius) and a PL quantum efficiency of $\sim 15\%$ at room temperature, was used for coating MF microspheres. The PL spectra from PM were recorded using a RENISHAW micro-PL system equipped with a positioning stage and an Ar⁺ laser ($\lambda = 488 \text{ nm}$, 1.0 mW power).

The small size of the MF spheres was dictated by specific requirements for the optimal excitation conditions like matching the laser wavelength to one of the WGM frequencies and achieving good correlation between the WGM and the laser linewidths. The wide separation of WGMs in spheres of this size allows us to avoid intricate band mixing in the PM. Moreover, the quality factor of $3 \mu\text{m}$ spherical particles forming the PM ($Q \sim 10^3$) provides better mode coupling than would be the case with larger microspheres⁹. We proceed from the assumption that a larger coupling parameter can be of crucial importance in the interaction between electromagnetic fields of two adjacent spheres, in order to experimentally study the fine spectral structure of coupled WGMs.

3. RESULTS AND DISCUSSION

To investigate the photon modes in interacting spherical microcavities we first measured PL spectra scanning a sample along the longitudinal axis of the PM (Scheme 1., Fig.1). The observed spectral structure originates from coupling of electronic transitions in NCs to the photon modes of the microsphere, with PL peaks corresponding to the

resonant frequencies of WGMs with transverse electric (TE) and transverse magnetic (TM) polarizations. In Fig. 1 one can clearly see a number of narrow peaks arising from TE and TM WGM of the individual spheres. The pronounced double structure, with the intensity distribution dependent on the excitation position, is just a result of superposition of the uncoupled WGMs of individual microspheres.

For spherically symmetric microparticles the extinction cross-section is derivable from Lorenz-Mie theory¹⁸:

$$C_{ext} = \frac{2\pi}{k^2} \text{Re} \sum_{n=1}^{\infty} (2n+1)(b_n(x, n_r) + a_n(x, n_r)), \quad (4)$$

where $x=2\pi r/\lambda$ is the size parameter, r is the radius of microsphere and the Mie scattering partial wave amplitudes $a_n(x, n_r)$ and $b_n(x, n_r)$ can be expressed in the form¹⁹

$$a_n(x, n_r) = \frac{A_n(x, n_r)}{A_n(x, n_r) + iC_n(x, n_r)} \quad \text{and} \quad b_n(x, n_r) = \frac{B_n(x, n_r)}{B_n(x, n_r) + iD_n(x, n_r)} \quad (5)$$

The pairs of functions A_n , C_n and B_n , D_n are determined by the Mie scattering problem²⁰ and resonance structure in scattering spectra can be expected as the real part of a_n or b_n reaches its maximum value of 1 and the imaginary part is passing through 0 from the positive to the negative side. In other words the resonances in the Mie scattering characteristics occur when $C_n = 0$ or $D_n = 0$, which gives the the following mathematical condition for a resonance²¹:

$$n_r \psi_n(x) \psi'_n(n_r x) - \psi_n(n_r x) \psi'_n(x) = 0 \quad (6)$$

or

$$n_r \chi_n(x) \psi'_n(n_r x) - \psi_n(n_r x) \chi'_n(x) = 0, \quad (7)$$

Note, that for given n and m , these equations have infinitely many solutions at discrete values of x for TM and TE WGMs respectively. The Riccati-Bessel functions of the first and second kind can be introduced as:

$$\psi_n(z) = z j_n(z) = \sqrt{\frac{\pi z}{2}} J_{n+\frac{1}{2}}(z); \quad \text{and} \quad \chi_n(z) = z y_n(z) = \sqrt{\frac{\pi z}{2}} Y_{n+\frac{1}{2}}(z); \quad (8)$$

where $J(z)$ and $Y(z)$ are the cylindrical Bessel functions of the first and second kind respectively. The use of Bessel functions for systems with cylindrical symmetry together with recurrence relation $\psi'_n(z) = -\frac{n}{z} \psi_n(z) + \psi_{n-1}(z)$

enables to reduce the equations (6)-(7) to the form convenient for practical calculation of position of WGMs. Thus for real refractive index conditions for TM and TE resonances can be taken in the form:

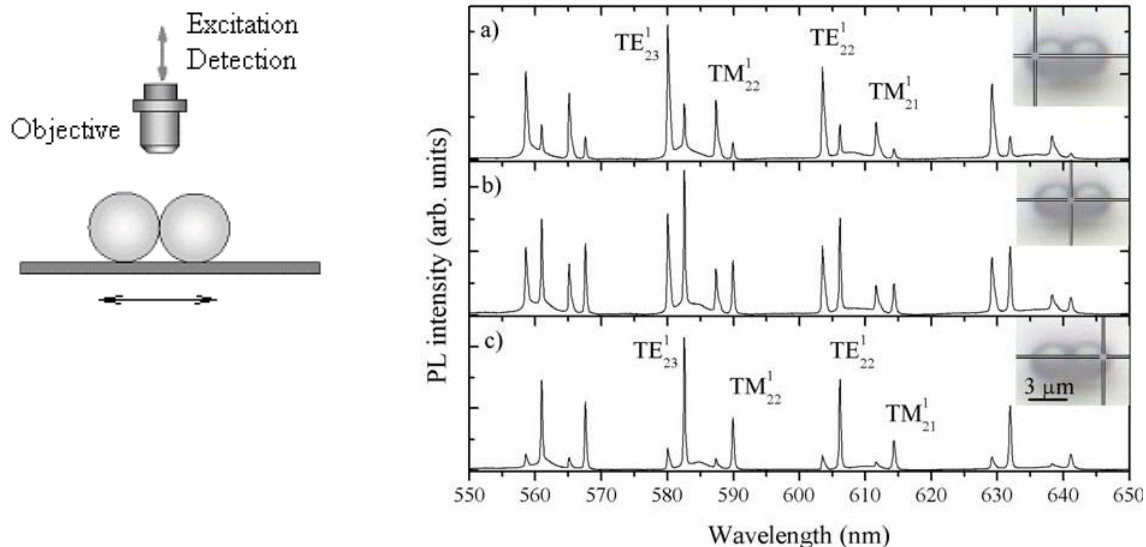
$$\frac{1}{x} \left(\frac{n}{n_r} - n_r n \right) J_{n+1/2}(x) J_{n+1/2}(n_r x) + n_r J_{n+1/2}(n_r x) J_{n-1/2}(x) - J_{n+1/2}(x) J_{n-1/2}(n_r x) = 0, \quad (9)$$

$$n_r Y_{n+1/2}(x) J_{n-1/2}(n_r x) - J_{n+1/2}(n_r x) Y_{n-1/2}(x) = 0. \quad (10)$$

These conditions are a transcendental equations, which can be solved for the size parameter x (position of a resonance) for given values of refractive index m and for given angular quantum number n . Thus comparing calculated results with the spectral positions of the WGM in the experimental PL spectra we can identify the indexes n and l for each mode and estimate the size of the sphere. The algorithm of the mode assignment can be as follows. 1) The resonant wavelengths corresponding WGM resonances λ_i^{exp} ($i = 1, 2, \dots, N$) are determined from a PL spectrum of single sphere. 2) We assume approximate value microsphere radius based on technical specification within distribution of the sizes. 3)

Theoretical resonance positions λ_i^{theor} are then calculated using eqs. (8) and (9). 4) Two lists are compared and for each value of λ_i^{exp} the closest value λ_i^{theor} is suggested and difference between them Δ_i is calculated. 5) Taking into account spectral resolution Δ the correlation $S = \frac{1}{N} \sum_1^N (1 + \Delta_i / \Delta)^{-1}$ is then maximized by adjusting only one fitting parameter, namely size of microsphere.

Figure 1 shows a result of WGM identification by this algorithm. The shift in position of the TE_{22}^1 , TE_{23}^1 , TM_{21}^1 and TM_{22}^1 modes of single microspheres, which can be clearly seen in Fig.1 is a result of a difference in the size of the two microspheres of only 11 nm as follows from our calculations based on Mie theory. When the excitation is provided at the contact point between microspheres (Fig.1,b), the excited and emitted light propagate efficiently along the longitudinal axis of the bisphere and the contribution of both microcavities can be clearly seen in the PL spectra at discrete WGM wavelengths. This contribution of PL signal from individual microcavities can be switched by changing excitation position from the left to the right rim of PM (Fig. 1, a, c). The above experiments have involved excitation beam incidence perpendicular to the longitudinal axis of the PM. For a given excitation and detection configuration we did not observe any new peaks in the PL spectra of the PM as compared to that of single spheres which is indication of weak intermode coupling. The lack of intermode coupling in this configuration (Scheme 1) is not surprising because the coupling between electro-magnetic fields of the spheres is expected to be minimum when the excitation and detection are set perpendicular to the axis of the PM⁹. However, the observed redistribution of intensity of components of WGM double structure clearly demonstrates the waveguiding of the light along the PM and the possibility of wavelength switching in the PM depending on excitation-detection geometry.



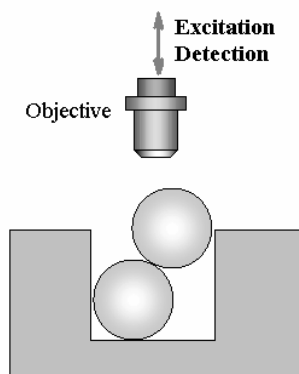
Scheme 1. Experimental geometry for scanning PM along the longitudinal axis.

Figure 1. PL spectra of a PM with excitation and detection at three different positions along its longitudinal axis. Insets: microscope images of the PM, with the cross-hairs indicating the excitation-detection position.

Theoretical considerations^{8, 10} show that the intersphere coupling is expected to be maximum for the pair of modes whose orbitals include the contact point between microspheres and lie in the same plane⁹. Taking this into account, the PL intensity of the coupled modes is anticipated to be maximum in the direction parallel to the PM axis and the signal from the coupled intersphere modes should be more pronounced in the parallel configuration than in the perpendicular one.

In order to control the alignment of the spheres constituting the PM we utilized a polystyrene substrate containing a three-dimensionally ordered array of pores of $\sim 5 \mu\text{m}$ in size prepared through a thermocapillary convection²². The ordered structures are formed by evaporating solutions of polystyrene in a volatile solvent, in the presence of moisture with forced airflow across the solution surface. A hexagonally packed array of holes (microwells) of 3-5 μm

depth then forms on the surface of the polymer (Fig.2). Only one pair of the 3 μm microspheres can be accommodated within each microwell, and the axis of the PM is close to the surface normal (Scheme 2).



Scheme 2. Experimental geometry for off-axis or parallel excitation and detection geometry.

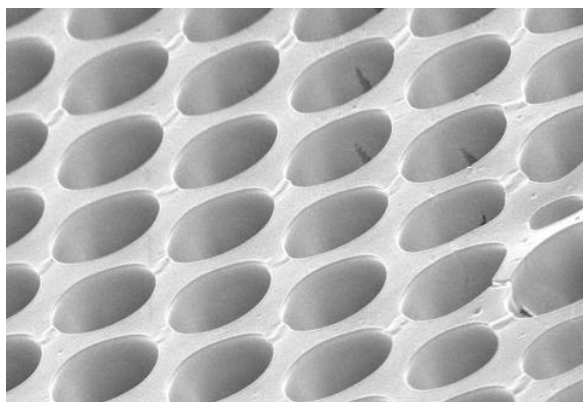


Figure 2. Polystyrene films with hexagonally ordered arrays of air holes with diameters of about 5 μm .

Figure 3 shows the PL spectrum of a PM accommodated in a microwell and the spectra of the individual microspheres prior to being manipulated into the microwell. Pronounced doublet structure in the spectrum presented in Fig. 3(a), coinciding with the spectral positions of WGM of individual microspheres (Fig. 3(b)), is a result of the overlap of uncoupled modes of the two spheres. The shift in position of the WGM, which can be clearly seen in Figs. 3(a) and 3(b), is a result of a difference in the size of the two microspheres of only 25 nm. However, in contrast to the spectrum in Fig. 3(b), this doublet is accompanied by two extra relatively broad peaks (indicated by arrows in Fig. 3(a)), which are indicative of strong mode coupling in the PM.

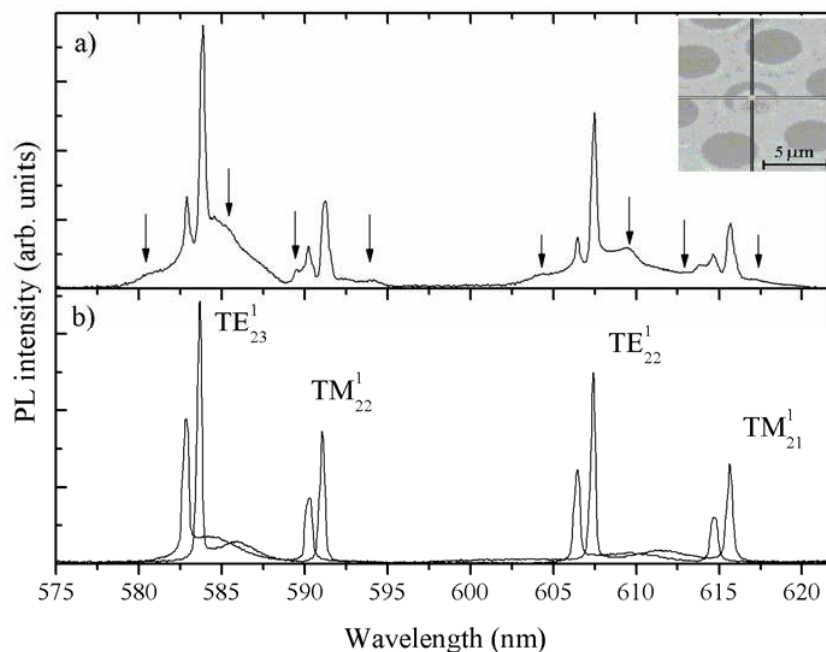


Figure 3. (a) The PL spectrum of a PM accommodated in a microwell. Arrows indicate the coupled modes. (b) The PL spectra of noninteracting microspheres. Inset: microscope image of the PM in the microwell. The dark cross indicates the excitation position.

The appearance of these two satellites can be interpreted as a result of the formation of BN and ABN orbitals in the PM⁷ with the ABN peak observed at lower wavelength than the BN one. In terms of cooperative scattering theory⁸, the observed satellites originate from the removal of the WGM degeneracy with respect to the m index. The line shape of the satellite lines reflects the energy distribution among the coupled modes, because modes with different combinations of m can contribute to the PL spectra. The observation of a broader ABN peak, relative to the BN peak, reflects the decrease of the quality factor of the PM when compared with that of a single sphere presumably due to the interaction with more dissipative modes of lower l ⁹. The deconvolution of the lineshape of resonances belonging to TE_{22}^1 WGM (Figure 3a,b) using Lorentian functions show that the quality factors of PM peaks are ~ 8 times smaller than Q-factors of noninteracting microspheres.

The high PL efficiency of CdTe NCs and the coupling of electronic transitions of NCs to the resonances of the PM allows us to detect the BN and ABN branches in a wide spectral region from 525 to 725 nm, corresponding to the full spectral range of the NC emission. This allows the possibility of estimating the magnitude of spectral spacing between bonding and antibonding branches as a function of angular mode number for comparison with theoretical calculations (Fig.4).

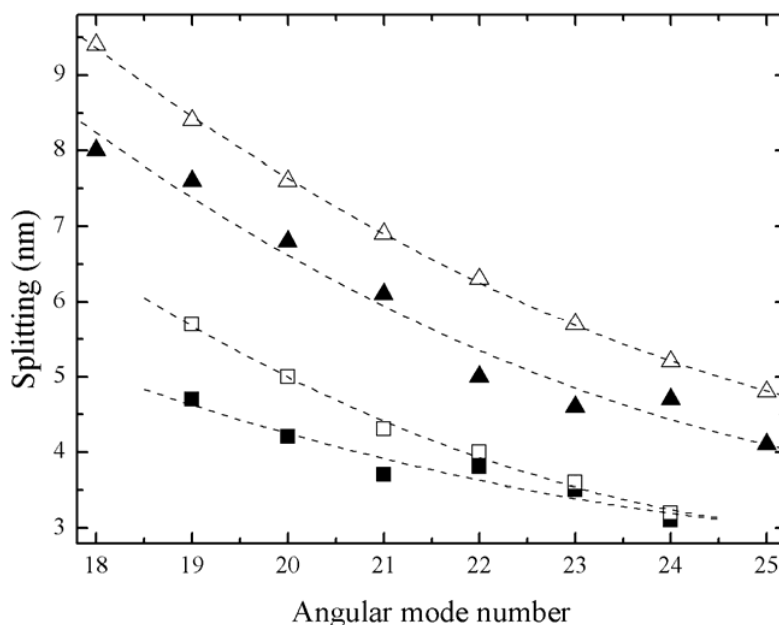


Figure 4. Magnitude of spectral spacing between bonding and antibonding branches estimated from micro-PL spectra of PM as a function of angular mode number for TE (solid up-triangles) and TM (solid squares) modes. Open symbols show the result of theoretical calculation. The dashed lines are a guide to the eye.

In order to calculate the BN/ABN splitting we adopted the maximum term approximation (MTA) of the single-mode tight-binding (SMTB) method¹⁰. For simplicity in our simulations we have considered only the case of identical spheres with WGM having radial quantum number $l = 1$, taking into account only the interaction between WGM of the same n .

For a given WGM polarization, the value of the splitting between the BN and ABN modes of PM can be obtained from

$$\Delta x_m = 2\Gamma_n |A_{n,m}(x_n)| \quad (11)$$

where $x_n = 2\pi R/\lambda_n$ is the size parameter of a resonance with mode number n , λ_n is the corresponding resonant wavelength, Γ_n is the width of mode n in a single sphere, which can be calculated within Mie-theory. For a given radius of single spheres a , the coefficient $A_{n,m}(x_n)$ can be calculated as:

$$A_{n,m}(x_n) = -2n(-1)^{n+m} h_{2n}^{(1)}(k_0 Z) \times \sqrt{\frac{2n}{\pi(n+m)(n-m)}} \times \frac{n^{2n}}{(n+m)^{n+m}(n-m)^{n-m}}, \quad (12)$$

where $k_0 = x_n/a$ is the wavenumber, Z is the characteristic length¹⁰, and the spherical Hankel function of the first can be estimated from

$$h_{2n}^{(1)}(k_0 Z) \cong -i \frac{\exp[(2n+1/2)(\alpha - \tanh\alpha)]}{(n+1/2)\sqrt{\sec h\alpha \tanh\alpha}}, \quad (13)$$

Here, α is defined by $k_0 Z = (2n+1/2)/\cosh \alpha$. The positions of the m -resonances were estimated as $x_m = x \pm \Delta x_m/2$.

Applying these equations we can see that the modeled data, presented in Fig. 4 displays behavior which is in good agreement with experimental data. From micro-PL spectra of the PM we have found a decrease in splitting between BN and ABN branches value with increasing n , with TE modes splitting being always higher than that of TM modes. Remarkably, the theory predicts the same behavior. Although the maximum term approximation overestimates the splitting (Fig. 4), it is clear that this approach provides a good qualitative guide for the analysis of a variety of phenomena observed in PM.

It is worth noting that observed disagreement between calculated and experimental values for the PM mode splitting may also have its origin in the fact that the magnitudes of splitting were estimated as a spectral distance between maxima of ABN and BN peaks. However detailed consideration of coherent mode coupling in PMs using SMTB approach shows that in fact the BN and ABN branches consist of a number of very narrow peaks, which are due to the presence of $m \neq \pm 1$ components¹⁰. The total number of these sharp peaks, originating from a certain mode n , is governed by the actual degeneracy of the Mie resonances which in the approach of the normal mode concept is $n+1$. (Due to the dependence of the mode coupling on the orbital plane orientation, the interaction is limited between modes of only the same m , no degeneracy is removed between m and $-m$ and the new degeneracy of PM modes is now $n+1$ with m running from 0 to n)¹⁰. In our calculations within the SMTB model we estimated the splitting as a spectral distance between outer peaks assigned to modes with $m=\pm 1$, which do not correspond to the maximum of envelope of PM modes and therefore this may cause discrepancy between estimated and observed values of the mode splitting. The fine structure of BN and ABN modes of PM is of great interest to the experimentalists because it suggests a manifold of applications particularly in the field of information processing¹⁶. However in order to observe this phenomenon two conditions should be met. First, the spacing between m -resonances forming BN and ABN modes of PM strongly depends on the angle of incidence of electromagnetic wave θ ^{9-10, 16}. For $\theta = 0^\circ$, when the incident light propagates parallel to the longitudinal axis of the PM, the incident wave can preferentially excite $m = \pm 1$ modes. The interaction with other dissipative modes of lower l causes the broadening of BN and ABN peaks and decreases the Q-factor relative to that of single spheres, as was observed in a number of papers^{8-9, 15-16}. In the perpendicular configuration, when $\theta = 90^\circ$, the intermode coupling and spacing between m -resonances are expected to be minimal. In other words, in the case of perpendicular orientation of the PM with respect to excitation, all m -modes would merge into broad BN and ABN features and their fine structure would not be easily recognized. A compromise can be reached only for off-axis excitation with $0^\circ < \theta < 90^\circ$. Secondly, in order to reveal the $m \neq \pm 1$ components interacting cavities should have not only similar size, but also similar Q-factors. If the resonances of the two cavities are of greatly different width, the coherent coupling is disturbed (at least to some extent) and therefore fine structure of PM modes cannot be detected.

One can see from the inset in Fig. 3a that excitation and detection in this case was provided at the centre of the upper sphere forming the PM i.e. along the PM axis. In order to reveal the $m \neq \pm 1$ components in the PL spectra of the MF/CdTe PM we have investigated WGM structure at different excitation (and detection) positions across the upper sphere.

Figure 5a shows micro-PL spectra of a PM formed by two almost identical microspheres (sample B) with sizes of 3.0168 μm and 3.0189 μm measured for excitation (and signal collection) at the right edge of the upper microsphere (inset, Fig.5a). The presented PL spectra clearly reveal major features unique to strong coherent coupling between the photonic states of the two microspheres forming the PM. One can clearly see a number of narrow peaks grouping on both sides arising from TE and TM resonances of the individual spheres and forming BN and ABN modes of PM. The origin of this fine structure of BN and ABN modes lies in the lifting of the mode degeneracy in the PM^{10, 16} and

therefore this multi peak structure would not be radically altered in scattering spectra or even in lasing, or stimulated Raman scattering in the same way that in a semiconductor laser the lasing does not alter the cavity resonances.

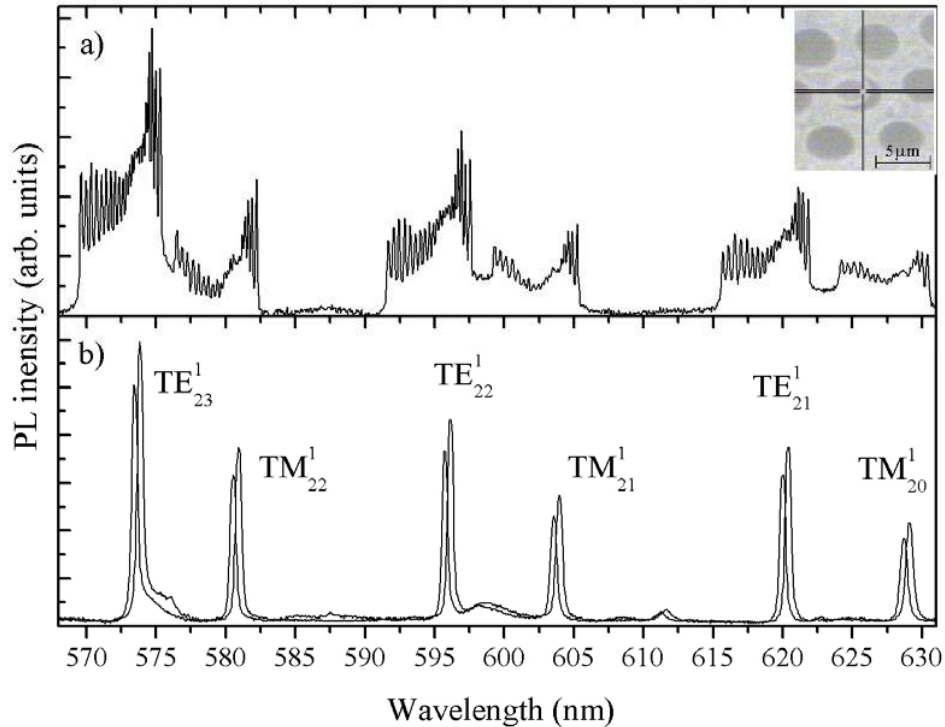


Figure 5. (a) PL spectra of the PM formed by two almost identical microspheres with off-axis excitation and detection. (b) The PL spectra of noninteracting microspheres. Inset: microscope image of the PM in the microwell. The dark cross indicates the excitation position.

It turns out that the number of experimentally resolved peaks increases with angular mode number both for TE and TM modes, and is very close to the n value, although never in excess of n . Indeed we have observed 19 peaks in the spectral region of TE_{20}^1 (19 peaks for TM_{20}^1), 21 peaks around the TE_{21}^1 resonance (20 peaks for TM_{21}^1), 21 peaks in the region of the TE_{22}^1 resonance (21 peaks for TM_{22}^1), 22 peaks for TE_{23}^1 (22 peaks for TM_{23}^1) and 23 peaks around the TE_{24}^1 resonance. Note that the spectral region occupied by a set of these narrow resonances is much wider than linewidth of WGMs of non-interacting microspheres (Fig.5).

The deconvolution of the lineshape of the m -resonances of the PM using Lorentian functions shows that m -resonances of the BN branch are always sharper than that of ABN one, providing a higher quality factor Q value for these modes and therefore a higher photon lifetime in the resonant modes τ . However, the most remarkable experimental fact is that the Q factor of the m -resonances in the spectra of the PM exceeds the Q value of single noninteracting microspheres, suggesting the respective modification of photon lifetime in PM (τ_{PM}) relative to photon storage time in single spherical microcavities before contact (τ_{SS}). These two facts along with estimated value of BN/ABN splitting (~ 5 -7 nm) implies the possibility for development of a new PM-based photonic device such as an optical delay line with controllable spectral and temporal tunability.

Figure 7 shows spectral distribution of the ratio between photon lifetime of m -modes of PM and photon storage time in WGMs of single sphere calculated from corresponding values of mode linewidth in the region of TE_{21}^1 and TM_{20}^1 resonances demonstrating a threefold photon storage enhancement. It is evident from Fig. 7, that the interaction between spherical microcavities results in periodic group delay spectra with peaks occurring at each of the m -resonant

frequencies with bigger delay times for higher m -values, which implies that the spectral components near these m -resonance spend more time traveling within the PM.

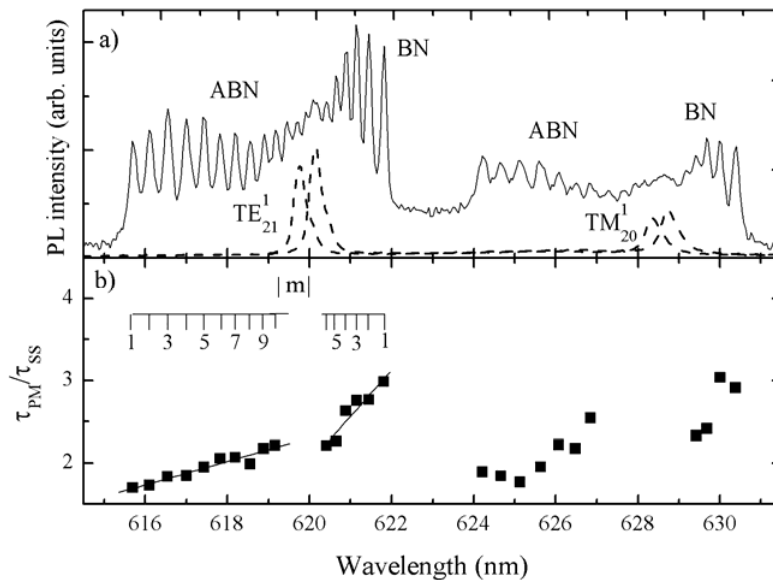


Figure 7. (a) PL spectra of the PM (solid line) formed by two almost identical melamine-formaldehyde microspheres with off-axis excitation and detection. Dashed lines show PL spectra of noninteracting microspheres. (b) Ratio between photon lifetime of the m -modes of photonic molecule and one of single spheres. The solid lines are result of linear fit.

The relative increase in photon storage time $R = \tau_{PM}/\tau_{SS}$ shows almost linear dependence on azimuthal number m for given mode number n (Fig. 7, 8), however it turned out that slope of this dependence ($\Delta R/\Delta m$) in its turn varies with n (Fig. 9, inset).

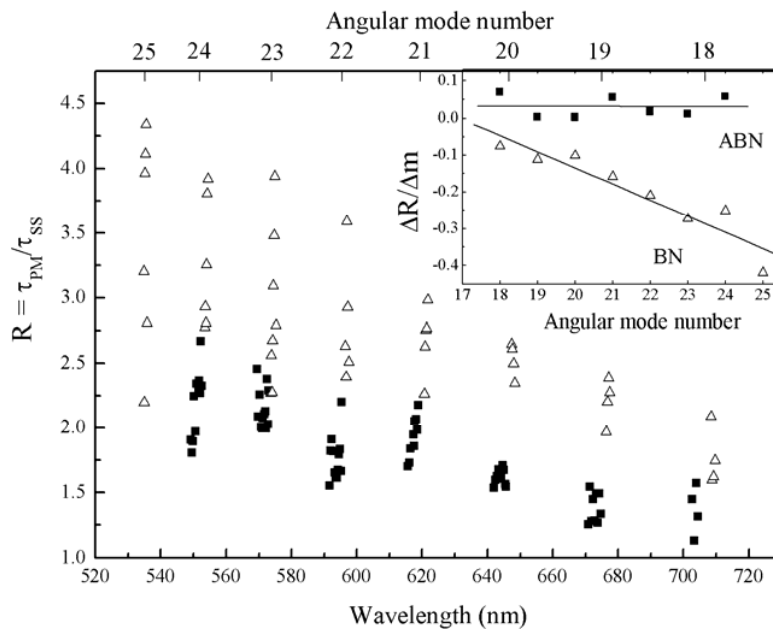


Figure 8. Ratio between photon lifetime of m -modes of photonic molecule (τ_{PM}) and photon storage time in WGM of single sphere (τ_{SS}). Inset shows n -dependence of slopes $\Delta R/\Delta m$ for BN and ABN modes. The solid lines are result of linear fit.

In our experiments we have investigated the strong coupling phenomenon in the spectra of the PM in the visible and near-IR spectral regions. In order to study the feasibility of the PM as an optical delay device, we have theoretically analysed the m -dependent structure for modes accommodated within the C band from 1525 – 1565 nm using MTA (Eq.11-13) and extrapolated the n -dependence of $\Delta R/\Delta m$ presented in Fig.8 into region of higher values of angular mode number.

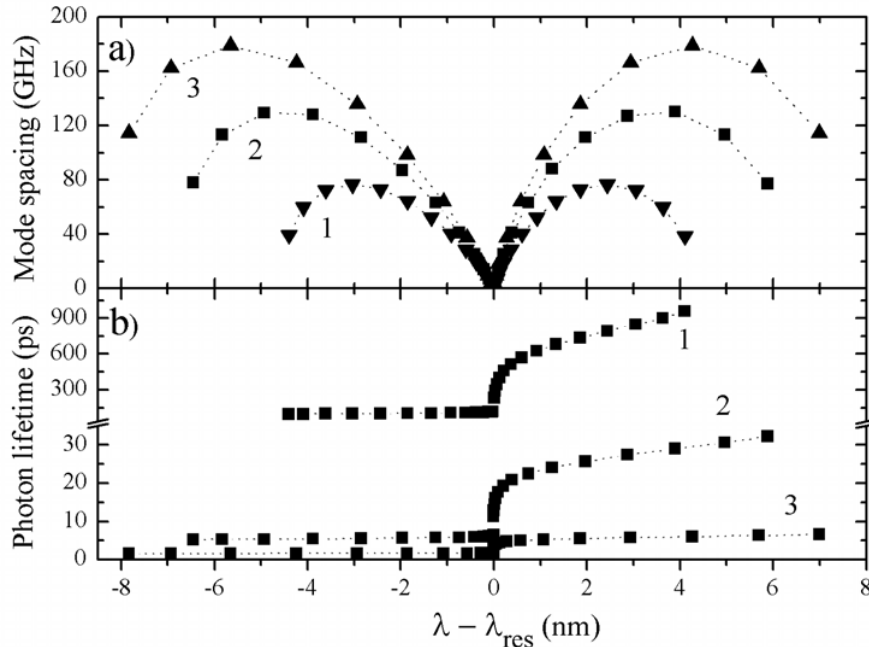


Figure 9. (a) Spacing between adjacent m -modes calculated in spectral region of the C-band for PMs formed by two identical SiO₂ microspheres with sizes of 10 μm (3), 12 μm (2) and 16 μm (1). (b) Calculated distribution of delay time between m -modes. The zero value on wavelength axis corresponds to the position of TE₂₄, TE₂₉ and TE₄₀ WGM of the single spherical microcavity.

Figure 9a shows the values of spacing between adjacent m -resonances calculated in spectral region of the C-band for PMs formed by two identical SiO₂ microspheres with sizes of 10 μm (3), 12 μm (2) and 16 μm (1). The standard for channel spacing in today's optical communication systems 40 GHz - 100 GHz²³ was taken into account deciding on microsphere sizes. In our model each m -resonance in the PM signal can be considered as a channel with spacing between channels depending on the size of the microspheres and the azimuthal mode number.

For 10-μm individual microcavities, only one TE mode ($\lambda_{24} = 1542.4$ nm) is found in the C-band. The maximum splitting between BN and ABN modes (i.e. for $m = 1$) of 15.7 nm was estimated in this case. Calculating the positions of resonances for a single 16-μm microsphere we have found 2 WGMs in the region of the C-band: TE₄₀($\lambda=1544.6$ nm) and TM₃₉($\lambda=1559.6$ nm). Maximum splitting in this case is much smaller: 8.8 nm for the TE₄₀ WGM, but accommodates a larger number of m -resonances. These results clearly demonstrate that two coupled WGM spherical microcavities can generate modal structure with controllable number of peaks distributed across the major communication band. In contrast to the case of single spheres, the bandpass of the proposed delay-line device can be as wide as the width of all communication bands and can be controlled by the size of spherical microcavities forming PM.

Figure 8b shows the estimated distribution of delay times between m -modes of the PM taking values of the photon storage time in single microspheres to be $\tau_{ss} = 55$ ps for TE₄₀¹, 3 ps for TE₂₉¹ and 0.9 ps for TE₂₄¹ as determined from our calculations based on the Mie-theory. We can anticipate that values of delay time, τ_{PM} , calculated in such a way are underestimated, because of the difference in Q-factor achievable for 3-μm spheres and 10-16μm ones. Recent experimental studies clearly show a strong dependence of the Q-factors of m -resonances (and therefore delay times) on quality factor of individual microspheres^{15,16}.

The observed non-uniformity of spacing between m -resonances (Fig.8, Fig.9a) provides a unique possibility to control the distribution of photon storage times in a particular spectral region (Fig.9b). For all sizes of interacting microspheres, we have obtained an initial increase of the spacing within the region of first three m -resonances followed by a dip and then an increase in spacing again (Fig. 8a). This theoretical result is in good qualitative agreement with experimental data¹⁶. For the smallest PM size, 18 m -modes are available within a window of 12 – 200 GHz. The smaller Q-factor of WGMs in the individual spheres forming the PM results in a relatively moderate increase in delay time – from 1.5 ps (obtained for ABN resonance with $m = 1$) up to 7 ps calculated for the corresponding m -resonance of BN mode of PM. For microsphere of larger size, the number of available modes increases. Indeed, for 12- μm coupled spheres 20 m -resonances fit into above indicated spacing window, with biggest value of intermode spacing \sim 130 GHz. Delay times distributed between these m -modes increase from 5 ps up to 32 ps. The upper limit of intermode spacing drops even more for PMs formed from 16- μm microspheres. In that case, 22 m -resonances with spacing between 12 and 76 GHz can be seen in Fig. 8a, providing discrete time delays found to be distributed between 94 ps and 960 ps. Remarkably, this maximum delay provided by two interacting microspheres of 16- μm size corresponds to a \sim 15 cm length of silica-glass waveguide delay line.

4. CONCLUSIONS

Using these results we can draw a few general conclusions. In the weak coupling regime interacting spherical microcavities provide the possibility of waveguiding and wavelength switching in the PM depending on the excitation-detection geometry. When coupling is strong, two spherical microcavities can generate modal structure with quality factors and photon lifetimes higher than that of single spheres. By analysing the m -dependent splitting of BN and ABN modes of PMs as a function of size, we have introduced a new strategy for the potential application of coupled spherical microcavities as a multi-channel, wavelength-tunable optical delay device. In the proposed scheme the bandpass width, mode spacings and delays depend only on the physical parameters of the interacting microcavities and have the desired property of being independent of the modulation frequency of the communication signal.

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